# 10. A Two-Stage Multi-Splitting Method for Non-Overlapping Domain Decomposition for Parabolic Equations 

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## Domain Decomposition

In domain decomposition for parabolic partial differential equations (PDE) several approaches have been developed- breaking the domain into multiple subdomains of either overlapping or non-overlapping type, or using algebraic type splittingscf. [CM94] for an overview. An important aspect is how to present the boundary conditions across interfaces or across common unknown points of subdomains, cf. [GS98, HT96, Tan92]. Towards parallelism, we divide the domain into subdomains with one grid point in common, adding an extra unknown at the interface to have effectively a non-overlapping decomposition.

In the present numerical method we have designed a one gridpoint overlap together with an extra equation in order to arrive at an effective multi-splitting approach. The transmission of data at the interface is through a discrete parametrized Robin boundary condition across interior interface points. A significant part of this report is the design and experimental study of optimizing boundary parameter coupled with particular choices of inner and outer splittings. We are interested here in extending some work of San and Tang [HT96] and Tang [Tan92] to parabolic problems. There is a parameter $\gamma$ that acts like a feedback gain across the artificial interfaces. The primary aspect of this article is to construct and demonstrate effective multi-splitting methods as depending on the interface boundary condition.

Consider the numerical solution of parabolic problems of form:

$$
\begin{equation*}
\frac{\partial u}{\partial t}=\frac{\partial^{2} u}{\partial x^{2}}+f(x, t), \quad x \in(0, L), t>0 \tag{1}
\end{equation*}
$$

subject to initial condition $u(x, 0)=g(x), x \in \Omega:=(0, L)$ and boundary conditions $u(0, t)=u(1, t)=0$.

With the usual $(\Delta x, \Delta t)$ mesh in $(x, t)\left(\Delta x=1 / M, x_{i}=i \Delta x, i=0, \cdots, M, t_{j}=\right.$ $j \Delta t, j=0, \cdots, N$ ), let $\bar{x}$ be a point of interface by which the domain $\Omega$ is decomposed into two subdomains $\Omega_{1}=\{x \in \Omega: x<\bar{x}+\Delta x\}$ and $\Omega_{2}=\{x \in \Omega: x>\bar{x}-\Delta x\}$, where $\bar{x}=m \Delta x$, for some $m$ with $1<m<M-1$. We are utilizing just two subdomains to address the essential issues. Solutions $u_{1}$ and $u_{2}$ are restricted versions of $u$ over the domains.

[^0]We introduce a boundary condition at $\bar{x}$ as follows

$$
\begin{equation*}
\alpha u_{1}+\left.(1-\alpha) \frac{\partial u_{1}}{\partial n}\right|_{\bar{x}-}=\alpha u_{2}+\left.(1-\alpha) \frac{\partial u_{2}}{\partial n}\right|_{\bar{x}+} \quad \alpha \in[0,1] . \tag{2}
\end{equation*}
$$

Then the global problem (1) is split over $\Omega_{1}$ and $\Omega_{2}$ in a natural fashion.
One can employ the method of lines with (1) through a second order central difference approximation [Smi85]: $\left.\frac{\partial^{2} u}{\partial x^{2}}\right|_{\left(x_{i}, t_{j}\right)}=\Delta x^{-2}\left(u_{i-1, j}-2 u_{i, j}+u_{i+1, j}\right)+O\left(\Delta x^{2}\right)$. The normal derivative in the boundary condition (2) at $(\bar{x}, t)=\left(x_{m}, t\right)$ is approximated by forward or backward differences for $\Omega_{1}$, and $\Omega_{2}$, respectively, as follows:

$$
\begin{align*}
& \left.\alpha u_{1}\right|_{m, j}+(1-\alpha) \frac{1}{\Delta x}\left(\left.u_{1}\right|_{m+1, j}-\left.u_{1}\right|_{m, j}\right)= \\
& \left.\alpha u_{2}\right|_{m, j}+(1-\alpha) \frac{1}{\Delta x}\left(\left.u_{2}\right|_{m+1, j}-\left.u_{2}\right|_{m, j}\right), \quad \bar{x}=x_{m} \in \Omega_{1},  \tag{3}\\
& \left.\alpha u_{2}\right|_{m, j}+(1-\alpha) \frac{1}{\Delta x}\left(\left.u_{2}\right|_{m, j}-\left.u_{2}\right|_{m-1, j}\right)= \\
& \left.\alpha u_{1}\right|_{m, j}+(1-\alpha) \frac{1}{\Delta x}\left(\left.u_{1}\right|_{m, j}-\left.u_{1}\right|_{m-1, j}\right), \quad \bar{x}=x_{m} \in \Omega_{1} . \tag{4}
\end{align*}
$$

In equation (3) $u_{2}$ is considered as given (known), $\left.u_{1}\right|_{m, j}$ is an unknown, and $\left.u_{1}\right|_{m+1, j}$ is a fictitious value for which this equation provides substitution. The situation is similar for equation (4). Re-write these equations by introducing a convenient parameter $\gamma:=(1-\alpha(1+\Delta x)) /(1-\alpha)$ as follows:

$$
\begin{align*}
& \left.u_{1}\right|_{m+1, j}-\left.u_{2}\right|_{m+1, j}=\gamma\left(\left.u_{1}\right|_{m, j}-\left.u_{2}\right|_{m, j}\right),  \tag{5}\\
& \left.u_{2}\right|_{m-1, j}-\left.u_{1}\right|_{m-1, j}=\gamma\left(\left.u_{2}\right|_{m, j}-\left.u_{1}\right|_{m, j}\right) . \tag{6}
\end{align*}
$$

In equations (5) and (6) the parametrized discrete Robin boundary condition at the matrix level (below) amounts to a kind of error feedback where $\gamma$ is the gain. If $\gamma$ is first chosen, then $\alpha$ in (3) and (4) becomes $\alpha=(\gamma-1) /(\gamma-(1+\Delta x))$. We note the particular choice of $\gamma=0(\alpha=1 /(1+\Delta x))$, giving a variation on the SAM Dirichlet condition in which an extra column has been inserted to slide the entries over; we call this the Sliding Dirichlet Condition.

By substituting (3) and (4) in (2), a couple of first order systems of differential for $u_{1}$ and $u_{2}$ arise:

$$
\begin{gather*}
\frac{d u_{1}}{d t}=B_{1} u_{1}+\left[\begin{array}{c}
0 \\
\vdots \\
0 \\
\Delta x^{-2}\left(u_{2, m+1}-\gamma u_{2, m}\right)
\end{array}\right],  \tag{7}\\
\frac{d u_{2}}{d t}=B_{2} u_{2}+\left[\begin{array}{c}
\Delta x^{-2}\left(u_{1, m-1}-\gamma u_{1, m}\right) \\
0 \\
\vdots \\
0
\end{array}\right], \tag{8}
\end{gather*}
$$

where $B_{1} \in \mathbf{R}^{m \times m}$, and $B_{2} \in \mathbf{R}^{M-m \times M-m}$.
The $i$-th row of $B_{1}$ has form $\Delta x^{-2}[0, \cdots, 0,1,-2,1,0, \cdots, 0], i=1, \cdots, m-1$, while the $m$-th row of $B_{1}$ is $\Delta x^{-2}[0, \cdots, 1,-2+\gamma]$. For $B_{2}$ we retain the numbering from the spatial grid, so the $i$-th row is $\Delta x^{-2}[0, \cdots, 0,1,-2,1,0, \cdots], i=m+$ $1, \cdots, M-1$, and the $m$-th row (the first row in actuality) is $\Delta x^{-2}[-2+\gamma, 1,0, \ldots, 0]$. Then we assemble the matrices $B_{1}$ and $B_{2}$ into a block matrix of coefficients $A$ to represent the global system of ordinary differential equations. There are now $M$ unknowns due to the extra unknown $u_{2, m}$ (formerly, there were $M-1$ unknowns). Setting $u=\left[u_{1}, u_{2}\right]^{T}$, we arrive at:

$$
\begin{equation*}
\frac{d u}{d t}=B u . \tag{9}
\end{equation*}
$$

Here, $B \in \mathbf{R}^{M \times M}$ is given by:

$$
B=\Delta x^{-2}\left[\begin{array}{cccccccc}
-2 & 1 & & & & & &  \tag{10}\\
& \ddots & \ddots & \ddots & & & & \\
& 1 & -2 & 1 & & & & \\
& & 1 & -2+\gamma & -\gamma & 1 & & \\
& & 1 & -\gamma & -2+\gamma & 1 & & \\
& & & & 1 & -2 & 1 & \\
& & & & & \ddots & \ddots & \ddots \\
& & & & & & 1 & -2
\end{array}\right] .
$$

The exact solution of the semi-discrete system (9) satisfies the following two-term recurrence relation, [Smi85, Var62]:

$$
\begin{equation*}
u(t+\Delta t)=e^{\Delta t B} u(t) \tag{11}
\end{equation*}
$$

Several algorithms for the numerical solution of (11) can be generated through an approximation to the exponential $e^{\Delta t B}$; in particular, we shall use rational functions via implicit Padé approximations, focussing on $(1,0)$ and $(1,1)$ Padé schemes, $c f$. [Smi85].

The $(1,0)$ and $(1,1)$ Padé approximations are given (respectively) by

$$
\begin{equation*}
e^{\Delta t B}=(I-\Delta t B)^{-1}+O(\Delta t) \&=(I-0.5 \Delta t B)^{-1}(I+0.5 \Delta t B)+O\left(\Delta t^{2}\right) . \tag{12}
\end{equation*}
$$

For each scheme of (12) the recurrence relation for $u(t+\Delta t)$ in (11) gives the following linear system to solve:

$$
\begin{equation*}
(I-\sigma B) u(t+\Delta t)=R_{m}(\Delta t B) u(t) \tag{13}
\end{equation*}
$$

where (for $(1,0)$ and $(1,1)$, respectively)

$$
R_{1}(\Delta t B)=I, \sigma=\Delta t, \& R_{2}(\Delta t B)=I+\sigma B, \quad \sigma=0.5 \Delta t
$$

Let $A=I-\sigma B$ and $\tau=\sigma \Delta x^{-2}$. Then $A$ is illustrated as follows:

$$
A=I-\sigma B=\left[\begin{array}{ll}
A_{11} & A_{12}  \tag{14}\\
A_{21} & A_{22}
\end{array}\right]
$$

where

$$
\begin{aligned}
& A_{11}=\left[\begin{array}{ccccc}
1+2 \tau & -\tau & 0 & \cdots & 0 \\
-\tau & 1+2 \tau & -\tau & & \vdots \\
0 & \ddots & \ddots & \ddots & 0 \\
\vdots & & -\tau & 1+2 \tau & -\tau \\
0 & \cdots & 0 & -\tau & 1+(2-\gamma) \tau
\end{array}\right], \\
& A_{12}=\left[\begin{array}{ccccc}
0 & 0 & \cdots & 0 & 0 \\
\vdots & & & & \vdots \\
0 & 0 & \cdots & 0 & 0 \\
\gamma \tau & \tau & 0 & 0 & 0
\end{array}\right], \\
& A_{21}=\left[\begin{array}{ccccc}
0 & 0 & \cdots & \tau & \gamma \tau \\
0 & 0 & \cdots & 0 & 0 \\
\vdots & & & & \vdots \\
0 & 0 & 0 & 0 & 0
\end{array}\right], \\
& A_{22}=\left[\begin{array}{ccccc}
1+(2-\gamma) \tau & -\tau & 0 & \cdots & 0 \\
-\tau & 1+2 \tau & -\tau & & \vdots \\
0 & \ddots & \ddots & \ddots & 0 \\
\vdots & & -\tau & 1+2 \tau & -\tau \\
0 & \cdots & 0 & -\tau & 1+2 \tau
\end{array}\right] .
\end{aligned}
$$

Due to the implicit nature of the finite difference scheme, we desire to avoid undue restriction on the size of $\Delta t$. Thus, we assume only that $\Delta t / \Delta x \leq 1$, so that $\tau$ is at most $O\left(\Delta x^{-1}\right)$ and is expected to be of order $\Delta x^{-1}$. We desire to choose $\alpha$ (or $\gamma$ ) to assure that $A$ is an $M$-matrix or an $H$-matrix. It is easy to check that one must have $\gamma \leq 1$ for $\alpha \in[0,1]$; if $0<\gamma \leq 1$ and $2 \gamma \tau \leq 1$, then $A$ is an $H$-matrix; and if $\gamma \leq 0$ then $A$ is an $M$-matrix. Since $\tau$ depends on the mesh size, we want to reject the condition $2 \gamma \tau \leq 1$. Therefore, we are now concentrating on the situation where $\gamma \leq 0$, in which case $A$ is an $M$-matrix. This is easy to check upon consulting [Var62, p. 85].

## The Two-Stage Multi-splitting Algorithm

The main advantage in the proposed boundary condition $(3,4)$ at interfaces is that it leads to a partitioning of the global matrix $A$ in (14) into particularly beneficial submatrix blocks, which are amenable to solution of multi-splitting iterative type. The parallel multi-splitting method of O'leary and White [OW85] to solve a linear system $A u=b$ is defined by considering multi-splittings from the decomposition $A=$ $M_{k}-N_{k}, k=1, \cdots, K$, such that each $M_{k}$ is invertible; we can form an iterative method, as follows:

$$
\begin{equation*}
u^{i+1}=M_{k}^{-1} N_{k} u^{i}+M_{k}^{-1} b . \tag{15}
\end{equation*}
$$

The multi-splitting iterates in (15) can be computed concurrently by introducing the non-negative diagonal matrices $E_{k}, \sum_{k=1}^{K} E_{k}=I$, and using them as follows:

$$
\begin{equation*}
u_{i+1}=H u_{i}+c, \tag{16}
\end{equation*}
$$

where $H=\sum E_{k} M_{k}^{-1} N_{k}$ and $c=\sum E_{k} M_{k}^{-1} b$.
For the linear system of concern here the first stage of splitting is according to the subdomains $\Omega_{k}$ (cp. [FS97]):

$$
M_{i}=\left[\begin{array}{lllll}
D_{11} & & & &  \tag{17}\\
& \ddots & & & \\
& & A i i & & \\
& & & \ddots & \\
& & & & D_{K K}
\end{array}\right]
$$

where the matrices $D_{k k}$ are the diagonals of $A_{k k}$, respectively, or any dummy diagonal matrices to make $M_{k}, k=1, \cdots, K$ invertible. We take $N_{k}=M_{k}-A$.

To calculate $u^{n+1}$ from $u^{n}$ (where $u=\left[u_{1}, u_{2}, \cdots, u_{K}\right]^{T}$ ), use:

$$
\begin{align*}
M_{k} w_{k}^{i+1} & =N_{k} u^{i}+b, \quad k=1, \cdots, K,  \tag{18}\\
u^{n+1} & =\sum_{k=1}^{K} E_{k} w_{k}^{i+1} .
\end{align*}
$$

This is rewritten as a single splitting with $M=\operatorname{diag}\left(A_{i i}\right)$.
With the splitting above one could implement a nice parallel algorithm. There are a wide variety of iterative schemes to consider. As in [FS94], it is useful to examine the possibility of introducing a second state of multi-splitting methods, this time splitting the matrices $M_{k}$ according to their convenient algebraic struture [OW85].

The second stage of splitting will be considered for the matrices $M_{k}, k=1, \cdots, K$. For this article, we deal with three rather standard approaches for inner iteration: the Jacobi, Gauss-Seidel, and SOR. Each subdomain block has been constructed to be of form only slightly different from a standard type matrix, and it is natural to test these schemes to gain insight into the proposed boundary treatment.

For brevity, we shall only describe the procedure, assuming the standard version of each algorithm is well known. The outer iteration (also called the first stage) consists of splitting the matrix $A$ as diagrammed earlier. Then we split the submatrices $A_{i j}$ to the right hand side whenever $i \neq j$ and on processor $i$ we split all other $A_{j j}$ submatrices to the right, leaving only their diagonals.

The Parallel Jacobi (or PJacobi) version of the inner iteration (or second stage) is to iterate $s$ times on processor $i$ for unknowns corresponding to the $i^{t h}$ subdomain, splitting the upper and lower diagonal parts to the right hand side as well.

The Parallel Gauss-Seidel (or PGauss-Seidel) inner iteration is to iterate $s$ times on processor $i$ for unknowns corresponding to the $i^{t h}$ subdomain, splitting the upper diagonal parts within that block to the right hand side.

Parallel Successive Over-Relaxation (or PSOR) does $s$ iterations for $i^{\text {th }}$ subdomain unknowns of standard SOR (using $\omega$ as the parameter).

## Numerical Experiments

In this section we demonstrate the performance of the algorithms described earlier and give empirical support to claim of convergence.

Let's work with the following prototype parabolic (heat) equation:

$$
\begin{gathered}
u_{t}=u_{x x}+x(1-x) \pi \sin (\pi t)-2 \cos (\pi t), \quad x \in(0,1), t>0 \\
u(x, 0)=x(x-1), \quad x \in(0,1),
\end{gathered}
$$

with boundary conditions $u(0, t)=u(1, t)=0$. The exact solution is $x(x-1) \cos (\pi t)$. We will compute approximations to the solution at target time $t_{*}=1.0$. All iterations have a halting signal of $\left\|u^{n+1}-u^{n}\right\| \leq 0.5 \Delta t^{2} \Delta x\left\|u^{n}\right\|$, which means smaller step sizes require more iterations.

The first numerical experiment to report, see Figure 1, is with inner splittings of Jacobi type, different values of $K$ (the number of subdomains), two particular choices of $\gamma$, and $s=1$. The computations were done in FORTRAN 90 on a Penitum III (not in parallel). The case $\gamma=0.0$ is called the Sliding Dirichlet Condition because it looks like a Dirichlet condition, shifted over by one column. It arises from $\alpha=1 /(1+\Delta x)$ in the discrete parametrized Robin condition, so it is not the ordinary Schwarz Alternating Method Dirichlet condition. It is observed in Figure 1 that the Sliding Dirichlet Condition allows the ordinary Jacobi algorithm to be parallelized in an efficient manner, in that the spectral radius appears to be nearly identical over any number of domain splittings (the numbers shown are rounded). Thus, our Parallel Jacobi method (PJacobi) would run a factor of $K$ times faster on a parallel machine with $K$ processors. (Here, communication is neglected because it is relatively insignificant.)

Moreover, there are better choices of $\gamma$. As shown in Figure 1, $\gamma=-1.0$ speeds up the PJacobi multi-splitting scheme quite significantly. (The value $\gamma=-1.0$ is not the absolute optimal from experiment, but is nearly so, and this is chosen for convenience.) This corresponds to a discrete parametrized Robin boundary condition with $\alpha=2 /(2+\Delta x)$.

Using old values from neighboring domains and a single inner iteration ( $s=1$ ), we want to have a look at the Gauss-Seidel scheme on each subdomain. The Parallel Gauss-Seidel scheme performs best when $\gamma$ is zero (or slightly positive); see Figure 2.

## References

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| K | $\gamma$ | $\begin{aligned} & \text { Avg. No. Iter. } \\ & \text { Per Step } \\ & \Delta x=0.003125 \end{aligned}$ | $\begin{aligned} & \text { Avg. No. Iter. } \\ & \text { Per Step } \\ & \Delta x=0.0015625 \end{aligned}$ |
| :---: | :---: | :---: | :---: |
| 1 | $\begin{aligned} & \hline \hline \text { N/A } \\ & \text { N/A } \end{aligned}$ | 4407 | 9856 |
|  |  |  |  |
| 2 | $\begin{gathered} \hline \hline 0.0 \\ -1.0 \\ \hline \end{gathered}$ | 4407 | 9856 |
|  |  | 4058 | 9436 |
| 4 | $\begin{gathered} \hline \hline 0.0 \\ -1.0 \end{gathered}$ | 4407 | 9856 |
|  |  | 3097 | 8070 |
| 8 | $\begin{gathered} \hline \hline 0.0 \\ -1.0 \\ \hline \hline \end{gathered}$ | 4407 | 9856 |
|  |  | 1632 | 5164 |
| 16 | $\begin{gathered} \hline \hline 0.0 \\ -1.0 \\ \hline \end{gathered}$ | 4407 | 9856 |
|  |  | 1468 | 2935 |
| 32 | $\begin{gathered} \hline \hline 0.0 \\ -1.0 \end{gathered}$ | 4407 | 9856 |
|  |  | 1486 | 2953 |

Figure 1: Parallel Jacobi showing the effect of $\gamma$ and $K(s=1)$.

| $K$ | Avg. No. Iter. <br> Per Step <br> $\Delta x=0.003125$ | Avg. No. Iter. <br> Per Step <br> $\Delta x=0.0015625$ |
| :---: | :---: | :---: |
| 1 | 941 | 1895 |
| 2 | 948 | 1905 |
| 4 | 948 | 1905 |
| 8 | 950 | 1906 |
| 16 | 957 | 1912 |
| 32 | 971 | 1927 |

Figure 2: Parallel Gauss-Seidel showing the effect of $K$ with $\gamma=0.0$ and $s=1$.
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| $K$ | Avg. No. Iter. <br> Per Step ( $\omega)$ <br> $\Delta x=0.003125$ | Avg. No. Iter. <br> Per Step ( $\omega)$ <br> $\Delta x=0.0015625$ |
| :---: | :---: | :---: |
| 1 | $108(1.90)$ | $167(1.92)$ |
| 2 | $118(1.90)$ | $214(1.90)$ |
| 4 | $142(1.87)$ | $217(1.90)$ |
| 8 | $170(1.84)$ | $243(1.89)$ |
| 16 | $231(1.78)$ | $321(1.85)$ |
| 32 | $365(1.64)$ | $510(1.75)$ |

Figure 3: Parallel SOR (at near optimal $\omega$ ) showing the effect of $K$ with $\gamma=0.0$ and $s=1$.


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